Effect of Cu Substitution on Thermoelectric Properties of Ge Clathrates

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We have prepared polycrystalline n-type Ba₉CuₓGa₉₋ₓGe₁₋₂y (x = 0–5, y = 16 – 3x) clathrate compounds by arc melting and spark plasma sintering techniques and investigated the effect of Cu substitution for Ge on their thermoelectric properties. The Hall carrier concentration for Cu-substituted compounds is almost constant at the order of 10²⁰ cm⁻³, which is comparable to that for Ba₉Ga₅Ge₁₀ stoichiometric compounds. The Seebeck coefficient for these compounds is comparable to that for n-Ba₉Ga₅Ge₁₀₋₃y compounds. From the analysis of the properties, the effective mass of the conduction band is estimated to be about 1.4 m₀, which is equivalent to or slightly smaller than that of n-Ba₉Ga₅Ge₁₀₋₃y compounds. The Hall mobility increases as the Cu composition increases. Its temperature dependence obeys approximately T⁻¹/² dependence in the range of 80–300 K, indicating the dominance of the alloy disorder scattering. Two models are discussed to account for the reduction in the alloy disorder scattering by Cu substitution.

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1. Introduction

Semiconducting clathrate compounds with a general formula M₅X₄₆ (M = Sr, Ba; X = Si, Ge, Sn) have been extensively studied as promising candidates for thermoelectric applications because they combine glasslike thermal conductivity and relatively good electronic transport properties.¹–⁸ The unit cell of type-I clathrates contains two pentagonal dodecahedra and six hexakaidecahedra formed by X atoms, which are fully disordered on the three distinct framework sites. The large cages inside the framework can incorporate alkali-earth atoms and their anharmonic vibration inside the cages reduces the lattice thermal conductivity through an interaction of the guest atoms with a wide spectrum of low frequency phonons. On the other hand, these materials maintain relatively good electronic conduction through the covalent sp³ hybridized framework, which dominates the band structure and electronic transport properties of the compounds.

Cordier and Woll⁹ reported the synthesis of some of transition-metal intermetallic compounds based on Ba₉Si₄₆ and Ba₉Ge₄₆ clathrates. In these compounds, Si or Ge atoms at the crystallographic 6c sites are replaced by transition metals. Recently, we have reported the effects of substitution of transition metals (Pt¹⁰ and Pd¹¹) for Ge on the electronic and thermoelectric properties of Ba₉Ga₄₆Ge₃₀-based compounds. The important points obtained from the study are as follows. In these compounds, the transition metals preferentially occupy at the crystallographic 6c sites, and the Seebeck coefficient is greatly enhanced by substitution of these transition metals. Furthermore, p-type clathrate compounds with varying carrier concentrations can be successfully prepared by Pd substitution.

To improve the thermoelectric figure of merit by this transition metal substitution, we have systematically investigated the thermoelectric properties of transition-metal-substituted Ge and Si clathrate compounds. In this paper, we report an investigation of the electronic and thermoelectric properties of polycrystalline Ba₉CuₓGa₉₋ₓGe₁₋₂y (x = 0–5, y = 16 – 3x) compounds with varying Cu compositions in an effort to elucidate the effect of transition metal substitution in Ge clathrate compounds.

2. Experimental Details

Polycrystalline Ba₉CuₓGa₉₋ₓGe₁₋₂y (x = 0–5, y = 16 – 3x) compounds were prepared by using arc melting and spark plasma sintering techniques. The composition (x and y = 16 – 3x) of the compounds was settled by considering the charge compensation for Ba²⁺ ions by substituting Cu and Ga: Cu nominally acts as a “trivalent” acceptor and Ga acts as a “monovalent” acceptor. Ba ingot (3N), Cu grain (4N), Ga ingot (6N), and Ge ingot (4N) were used as the starting materials. Ba₉CuₓGa₉₋ₓGe₁₋₂y ingots were first prepared by arc melting constituent elements in an argon atmosphere. The obtained ingots were ground to fine powder (<90μm) and sintered at 988–1088 K and 30 MPa for 60 min in an argon atmosphere by the spark plasma sintering method. The density of each sample was measured by the Archimedean technique. Powder X-ray diffraction (XRD) and electron probe microanalysis (EPMA) of polished surface of samples were employed to check the crystal structure and the chemical composition of samples.

The electrical conductivity σ and the Seebeck coefficient S were measured in the temperature range of 300–900 K. Hall measurements were performed in the temperature range of 80–300 K by applying an external magnetic field of 1 T using the van der Pauw’s method. The Hall mobility µ was determined from the electrical conductivity σ and the Hall coefficient Rₓ using µ = σRₓ.
3. Results and Discussion

Figure 1 shows X-ray diffraction patterns for Ba$_{8}$Cu$_{2}$Ga$_{x}$Ge$_{46-x-y}$ (x = 0, 1, 3, 5; y = 16 - 3x) compounds. The X-ray diffraction patterns of all samples were indexed to a cubic unit cell corresponding to the clathrate structure (Cubic, space group $Pm\bar{3}m$). Although a weak diffraction peak due to unreacted Ge phase was observed for the Cu composition x = 5, the samples with x = 0–4 were found to be single phase materials.

Table 1 summarizes some room-temperature properties for Ba$_{8}$Cu$_{2}$Ga$_{x}$Ge$_{46-x-y}$ (x = 0–5; y = 16 – 3x) compounds. EPMA confirmed the homogeneity of the samples with x = 0–4 and the existence of a trace amount of unreacted Ge phase in the sample with x = 5. As shown in Table 1, it was also found from EPMA that the samples had almost the same compositions as the nominal ones that expected from the amounts of constituent elements in the starting mixtures. The analysis of XRD data showed that the lattice constant decreased linearly as the Cu composition x increased, indicating the substitution of Cu for Ge sites. The extrapolation to x = 6 from the experimental dependence agrees well with the value (1.06859 nm$^3$) for Ba$_{8}$Cu$_{6}$Ge$_{40}$. This change of the lattice constant can be explained in terms of the Vegard’s law. The Hall coefficient for all samples was found to be negative and almost independent of the temperature between 80–300 K, indicating dominant n-type conduction in nearly degenerate range. The Hall carrier concentration at room temperature of the samples was nearly constant at the order of 10$^{20}$ cm$^{-3}$. Due to the charge compensation for Ba ions, the carrier concentration for Cu-substituted samples is comparable to that for the Ba$_{8}$Ga$_{16}$Ge$_{30}$ stoichiometric compound.

Since the thermoelectric properties depend strongly on the carrier concentration, we need to isolate the effect of Cu substitution from the carrier concentration dependence of the samples with slightly different carrier concentrations. Thus, we compared the carrier concentration dependence of the room temperature electrical conductivity and Seebeck coefficient for Ba$_{8}$Cu$_{2}$Ga$_{x}$Ge$_{46-x-y}$ (x = 1–5; y = 16 – 3x) compounds to that for Ba$_{8}$Ga$_{16}$Ge$_{30}$ compounds, as shown in Figs. 2 and 3, respectively. It is evident from Fig. 2 that the electrical conductivity $\sigma$ is significantly enhanced by Cu substitution at similar carrier concentration as compared to n-Ba$_{8}$Ga$_{16}$Ge$_{30}$ compounds. On the other hand, the Seebeck coefficient for Ba$_{8}$Cu$_{2}$Ga$_{x}$Ge$_{46-x-y}$ (x = 1–5; y = 16 – 3x) compounds is comparable to or slightly smaller than that for n-Ba$_{8}$Ga$_{16}$Ge$_{30}$ compounds. We estimated the effective mass assuming a single parabolic band model with alloy disorder scattering as dominant scattering mechanism, as will be discussed later. We first calculated the reduced Fermi energy $\eta$ from the experimental Seebeck coefficient data using

$$S = -\frac{k_B}{e} \left[ \frac{2F_1(\eta)}{F_0(\eta)} - \eta \right],$$

(1)

where $k_B$ is Boltzman’s constant, $e$ is electron charge and $F_0$ is a Fermi integral of order x. The carrier concentration $n$ can be expressed as

$$n = 4\pi \left( \frac{2m^*k_B T}{\hbar^2} \right)^{1/2} F_1(\eta),$$

(2)

where $m^*$ is the effective mass, $T$ is the temperature, and $\hbar$ is Plank’s constant. Thus, the effective mass was calculated by using the calculated $\eta$ and the experimental carrier concent-

<table>
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<th>$x$</th>
<th>Nominal composition</th>
<th>EPMA composition</th>
<th>Density, $D$/g cm$^{-3}$</th>
<th>Lattice constant, $a_0$/nm</th>
<th>Carrier concentration, $n$/10$^{20}$ cm$^{-3}$</th>
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<td>Ba$<em>8$Ga$</em>{16}$Ge$_{30}$</td>
<td>Ba$<em>{27.68}$Ga$</em>{16.07}$Ge$_{29.45}$</td>
<td>5.782 (99.5%)</td>
<td>1.07874 ± 0.00024</td>
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<td>1</td>
<td>Ba$<em>8$Cu$</em>{1}$Ga$<em>{15}$Ge$</em>{31}$</td>
<td>Ba$<em>{27.68}$Ga$</em>{16.07}$Ge$_{29.45}$</td>
<td>5.792 (99.1%)</td>
<td>1.07667 ± 0.00016</td>
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<td>3</td>
<td>Ba$<em>8$Cu$</em>{3}$Ga$<em>{13}$Ge$</em>{34}$</td>
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<td>5.859 (99.5%)</td>
<td>1.07312 ± 0.00016</td>
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<td>4</td>
<td>Ba$<em>8$Cu$</em>{4}$Ga$<em>{12}$Ge$</em>{36}$</td>
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<td>5.820 (97.7%)</td>
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</table>
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Fig. 2 Electrical conductivity $\sigma$ at room temperature as a function of carrier concentration $n$ for $\text{Ba}_8\text{Cu}_x\text{Ga}_y\text{Ge}_{46-x,y}$ ($x = 1-5, y = 16 - 3x$) compounds.

Fig. 3 Seebeck coefficient $S$ at room temperature as a function of carrier concentration $n$ for $\text{Ba}_8\text{Cu}_x\text{Ga}_y\text{Ge}_{46-x,y}$ ($x = 1-5, y = 16 - 3x$) compounds.

Fig. 4 Hall mobility $\mu$ as a function of temperature for $\text{Ba}_8\text{Cu}_x\text{Ga}_y\text{Ge}_{46-x,y}$ ($x = 0-5, y = 16 - 3x$) compounds.

The transport properties in semiconductors are determined by the band structure and various scattering mechanisms. To elucidate the origin of the enhancement in the electrical conductivity, we investigated the carrier mobility and the scattering mechanism of $\text{Ba}_8\text{Cu}_x\text{Ga}_y\text{Ge}_{46-x,y}$ ($x = 0-5, y = 16 - 3x$) compounds. Figure 4 shows the Hall mobility as a function of temperature for $\text{Ba}_8\text{Cu}_x\text{Ga}_y\text{Ge}_{46-x,y}$ ($x = 0-5, y = 16 - 3x$) compounds. The Hall mobility for $\text{Ba}_8\text{Cu}_x\text{Ga}_y\text{Ge}_{46-x,y}$ ($x = 0-5, y = 16 - 3x$) compounds obeys approximately $\mu \propto T^{-1/2}$ behavior in the temperature range of 80–300 K, indicating that the alloy disorder scattering is the dominant scattering mechanism. Furthermore, it is interesting to note that the Hall mobility tends to increase with increasing Cu composition $x$. The enhancement of Hall mobility is also clearly identified from Fig. 5, which shows the Cu composition dependence of the room temperature Hall mobility for $\text{Ba}_8\text{Cu}_x\text{Ga}_y\text{Ge}_{46-x,y}$ ($x = 0-5, y = 16 - 3x$) compounds. The Hall mobility values for samples with high Cu compositions ($x = 4, 5$) are about two times larger than that for the $\text{Ba}_8\text{Ga}_{16}\text{Ge}_{30}$ ($x = 0$) sample. Since the carrier concentration for $\text{Ba}_8\text{Cu}_x\text{Ga}_y\text{Ge}_{46-x,y}$ ($x = 1-5, y = 16 - 3x$) compounds is nearly equivalent to that for $\text{Ba}_8\text{Ga}_{16}\text{Ge}_{30}$ ($x = 0$) compounds, the observed enhancement of the electrical conductivity $\sigma$ is attributed mainly to the increased Hall mobility.

We finally discuss possible models to account for the observed reduction of alloy disorder scattering by Cu substitution for $\text{Ba}_8\text{Cu}_x\text{Ga}_y\text{Ge}_{46-x,y}$ ($x = 0-5, y = 16 - 3x$) compounds. As reported by Cordier and Wolf, one of important crystallographic features in transition-metal clathrate compounds, including $\text{Ba}_8\text{Cu}_6\text{Ge}_{40}$ and $\text{Ba}_8\text{Cu}_6\text{Si}_{42}$, is the preferential occupation of transition metal atoms at the 6c sites. This feature was also observed for a series of $\text{Ba}_8\text{X}_6\text{Ga}_y\text{Ge}_{46-x,y}$ compounds with $X = \text{Pd}, \text{Pt}, \text{Ag},$ and $\text{Au}$. Similarly, if the Cu atoms preferentially occupy at the 6c site, the scattering of carriers may decrease due to the local ordering in the framework of the clathrate structure, which determines the electronic transport properties. Since the Ga composition decreases when the Cu composition increases, the total amount of atoms $(x + y = 16 - 2x)$
substituted for Ge decreases, resulting in the decrease in alloy disorder scattering. In this model, the electron mobility \( \mu_A \) can be expressed by Brook’s formula:\(^{15}\)

\[
\mu_A = \frac{3\alpha(1-\alpha)m^*5/2}{\sqrt{2\pi m^* h^4 N_0}} \left( \frac{1}{\Delta E} \right)^3 \sqrt{\frac{\kappa B T}{m^*}},
\]

where \( N_0 \) is the number of atoms per unit volume, \( m^* \) is a band-edge effective mass, \( \alpha \) is the fractional concentration of one of the species, and \( \Delta E \) is an energy parameter characterizing the alloy potential fluctuations. \( \text{Ba}_8\text{Cu}_y\text{Ga}_{1-y}\text{Ge}_{46-x} \) (\( x = 0-5, y = 16-3x \)) compound can be described as an alloy between \([\text{Ba}_8\text{Ge}_{46}]_a\) and \([\text{Ba}_8(\text{CuGa})_{16}]_b\). The dashed line in Fig. 5 is a fit of the experimental data to the alloy disorder scattering model by Eq. (3) with \( m^* = 1.4m_0 \) and \( \Delta E = 0.92\text{ eV} \). In \( \text{Si}_{1-x}\text{Ge}_x \) alloy, the energy parameters \( \Delta E \approx 0.5\text{ eV} \) and \( \approx 0.6\text{ eV} \) have been reported based on this model.\(^{15}\) As seen in Fig. 5, the composition dependence of the Hall mobility can be explained well in terms of alloy disorder scattering model. The relatively large disagreement for \( x = 5 \) may be due to the influence of an impurity (unreacted Ge) phase in the sample. We also estimated the relaxation time \( \tau \) from the relation \( \tau = \mu m^*/e \). Figure 6 shows the Cu composition dependence of the relaxation time \( \tau \) at room temperature for \( \text{Ba}_8\text{Cu}_y\text{Ga}_{1-y}\text{Ge}_{46-x} \) (\( x = 0-5, y = 16-3x \)) compounds. It turns out that the relaxation time \( \tau \) increases remarkably with increasing Cu composition. From these results, it is found that Cu substitution has a significant effect on the carrier scattering mechanism.

4. Conclusions

We have investigated the effect of Cu substitution for Ge on the thermoelectric properties of polycrystalline \( \text{Ba}_8\text{Cu}_y\text{Ga}_{1-y}\text{Ge}_{46-x} \) (\( x = 0-5, y = 16-3x \)) clathrate compounds prepared by arc melting and spark plasma sintering techniques. \( \text{Ba}_8\text{Cu}_y\text{Ga}_{1-y}\text{Ge}_{46-x} \) (\( x = 1-5, y = 16-3x \)) compounds have nearly degenerate \( n \)-type conduction, as is the case for \( \text{Ba}_4\text{Ge}_{10} \) stoichiometric compounds. The Hall mobility is significantly enhanced by substituting Cu for Ge while remaining relatively large Seebeck coefficient due to the effective mass of conduction band comparable to or slightly smaller than that for \( n-\text{Ba}_8\text{Ga}_{46} \) compounds. The temperature dependence of the Hall mobility indicates that the dominant scattering mechanism of conduction electrons is alloy disorder scattering in the temperature range of 80–300 K. We discussed two models to account for the reduction of carrier scattering by Cu substitution. One is the effect of local ordering in the framework structure due to the preferential occupation of Cu at the crystallographic (6c) site. The other is the effect of large decrease in the amount of atoms substituted for Ge by using Cu with the valence higher than Ga for charge compensation for Ba ions. The observed dependence of Hall mobility on the Cu composition can be explained in terms of the alloy disorder scattering model considering the later case. Although the clathrate compounds have interesting, favorable physical properties based on a phonon glass and electron crystal concept, the carrier mobility in \( n-\text{Ba}_8\text{Ga}_{46} \) compounds is still smaller than that in state-of-the-art thermoelectric materials. Our result, therefore, suggests that the transition metal substitution is an approach to optimize the electronic properties of clathrate compounds for thermoelectric applications.

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